

Numerical Methods for Parameter Identification in a Convection-Diffusion Equation

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Abstract

We use a recently developed Sinc-Galerkin method for the solution of non-self-adjoint equations to solve a parameter identification problem arising from the one-dimensional convection-diffusion equation. This method is well suited for unbounded domains and certain singularities in the coefficients, as we illustrate by several examples. Practical aspects of implementation of this method are considered in detail.

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1 Introduction

Sinc methods for the numerical solution of ordinary differential equations have been found to be a very effective technique, particularly for problems with singular solutions and those on unbounded domains. In this paper we consider an inverse problem associated with the direct problem of solving a steady state convection-diffusion system in one space dimension given by

$$\begin{aligned} Lc \equiv -(D(x)c_x)_x + (v(x)c)_x + \lambda(x)c &= f(x), \quad x > 0 \\ Ac(0) + Bc_x(0) &= G, \quad B \neq 0 \\ c(\infty) &= 0. \end{aligned} \tag{1}$$

The inverse problem is that of determining the diffusion coefficient $D(x)$ from measurements of the solution $c(x)$. This direct problem is singular since it is defined on a semi-infinite interval. Moreover, we allow certain singularities in the coefficients $v(x), \lambda(x)$ at the endpoints. We make assumptions about the coefficients and the solution as well. In particular, we assume that the dispersion coefficient $D(x)$ is continuous, bounded and tends to a constant as $x \rightarrow \infty$. These conditions ensure well-posedness of the direct problem. Additional conditions are required for Sinc methods to work and we discuss these in Section 3.

There are many Sinc approaches to solving second order differential equations (see, e.g., [4, 6] for comprehensive surveys.) These typically require a unit coefficient for the second order term of the differential equation, which would necessitate differentiation of the diffusion coefficient $D(x)$. Methods that avoid this differentiation were developed in [2] (see also [4, p. 188]) for selfadjoint problems on a finite interval with Dirichlet boundary conditions. The direct solver that we employ was developed in [5] to also avoid differentiation of the diffusion coefficient $D(x)$ and to handle the case of a non-selfadjoint problem with mixed boundary conditions as well. We shall show that when properly tuned, it yields a robust and effective tool for identification of $D(x)$. Electronic copies of the programs used in this paper are available upon request from the author. These programs were written in the language Octave (see [1]) but are easily ported to MATLAB.

2 Algorithm for the direct problem

2.1 Sinc background

Recall that the *sinc function* is defined for all $z \in \mathbb{C}$ by

$$\operatorname{sinc}(z) = \begin{cases} \frac{\sin(\pi z)}{\pi z} & z \neq 0 \\ 1 & z = 0. \end{cases}$$

Let h be a positive constant. We will denote the *sinc basis functions* by

$$S(k, h)(x) = \operatorname{sinc}\left(\frac{x - kh}{h}\right) \quad k \in \mathbb{Z}, \quad -\infty < x < \infty.$$

The sinc differentiation operator matrix of dimension $m + 1$, is given by $I^{(1)} = [\delta_{jk}^{(1)}]_{j,k=0}^m$, where $\delta_{jk}^{(1)} \equiv hS'(j, h)(x)|_{x=x_k}$ and similarly $I = I^{(0)} = [\delta_{jk}] = [S(k, h)(jh)]$.

Let ϕ be a conformal map which maps the horizontal strip of radius d about the x -axis onto the domain \mathcal{D} with boundary Γ and inverse map ψ and $F(z)$ an element of

$$B(\mathcal{D}) = \left\{ F(z) \mid F \text{ is analytic on } D, \int_{\Gamma} |F(w)dw| < \infty \right\}$$

Assume $F(x)$ is real valued on $(0, \infty)$, γ is analytic, $\phi'F/\gamma \in B(\mathcal{D})$ and suppose there are positive constants α, β and C so that

$$\left| \frac{F(\xi)}{\gamma(\xi)} \right| \leq C \quad \begin{cases} \exp(-\alpha|\phi(\xi)|), & \xi \in \Gamma_a \\ \exp(-\beta|\phi(\xi)|), & \xi \in \Gamma_b \end{cases}$$

where $\Gamma_a \equiv \{\xi \in \Gamma : \phi(\xi) = x \in (-\infty, 0)\}$ and $\Gamma_b \equiv \{\xi \in \Gamma : \phi(\xi) = x \in [0, \infty)\}$.

Theorem 2.1. *Assume notation as above and make parameter selections*

$$N = \lceil \frac{\alpha}{\beta} M \rceil \quad \text{and} \quad h = \left(\frac{\pi d}{\alpha M} \right)^{\frac{1}{2}} \leq \frac{2\pi d}{\ln(2)}.$$

Set $w_k = \psi(kh)$. Then there exists constants K, L independent of M and $m = 0, 1, \dots, n$, such that for all $\xi \in \Gamma$

$$\left\| \frac{d^m}{d\xi^m} F(\xi) - \sum_{k=-M}^N \frac{F(w_k)}{g(w_k)} \frac{d^m}{d\xi^m} \gamma(\xi) S(k, h) \circ \phi(\xi) \right\|_{\infty} \leq KM^{(m+1)/2} e^{-\sqrt{\pi d \alpha M}}$$

and

$$\left| \int_a^b (vw[S(j, h) \circ \phi])(x)dx - h \frac{vw}{\phi'}(x_j) \right| \leq LM e^{-\sqrt{\pi d \alpha M}}.$$

Proofs of these facts can be found in the Stenger text [6] or the Lund and Bowers text [4]. This fact exhibits the exponential accuracy of correctly tuned sinc approximations. It also underscores the need to understand the asymptotic nature of the function to be approximated. Extensive discussions of these matters can be found in the texts cited.

2.2 Variational and numerical variational forms

Following [5], we multiply the differential equation of system (1) by suitable test functions u which vanish at the endpoints, and use integration by parts to obtain

$$\int_0^\infty (Dc_x u_x - vcu_x + \lambda cu)dx = \int_0^\infty f(x)u dx.$$

We can express this variational form as $T(c, u) = R(u)$. Now define

$$\tilde{c}(x) \equiv c(x) - c'(0)q_0(x) - c(0)q_1(x)$$

for suitable cardinal functions $q_0(x)$, $q_1(x)$ satisfying

$$q_0(0) = 0, q_1(0) = 1, q_0'(0) = 1, q_1'(0) = 0$$

and we obtain that $\tilde{c}(0) = 0$ and $\tilde{c}'(0) = 0$. Sinc approximation theory can be applied to this function and it yields the *numerical variational form*

$$T(\tilde{c}, u) + c_x(0)T(q_0, u) + c(0)T(q_1, u) = R(u).$$

For a Sinc-Galerkin approximation we choose test functions $u_j(x) = \gamma(x)S(j, h) \circ \phi(x)$ where $\gamma(x)$ is a weight function and $\phi(x)$ is a conformal map from $(0, \infty)$ to $(-\infty, \infty)$, such as $\phi(x) = \ln(x)$. We assume that the coefficients of the problem (1) and the choice of cardinal functions q_0 and q_1 are such that the solution $\tilde{c}(x)$ satisfies the restrictions of Theorem 2.1. Thus we can approximate \tilde{c} exponentially well by

$$\tilde{c}_m(x) = \sum_{j=-M_x}^{N_x} \frac{d_j}{\gamma(w_j)} \gamma(x)S(j, h) \circ \phi(x)$$

where γ is an appropriately chosen so that $\tilde{c}'(x)$ is accurately approximated by the derivative of the cardinal sum and $m = M_x + N_x + 1$, the number of

sinc nodes. More importantly, this formula can be differentiated to obtain an exponentially accurate formula for the derivative of $\tilde{c}'(x)$. This formula can be inserted into the numerical variational form for the direct problem.

Let $d = [d_j]_{j=1}^m$ be the vector of sinc node values of $\tilde{c}_m(x)$. The discrete system corresponding to the numerical variational form can, after much algebraic manipulation, be expressed as a linear system with coefficient matrix M , namely

$$Md + c'_0 T_0 + c_0 T_1 = R_{dis},$$

where $c_0 = c(0)$ and $c'_0 = c_x(0)$. Let $w_s = M^{-1}R_{dis}$, $w_0 \equiv M^{-1}(T_0)$ and $w_1 \equiv M^{-1}T_1$ and we may write this equation in the form

$$d = w_s - c'_0 w_0 - c_0 w_1.$$

A scheme for computing c'_0 and c_0 can be derived by integrating the differential equation of the problem from 0 to ∞ to obtain

$$\int_0^\infty (f(x) - \lambda(x)c)dx = D(0)c_x(0) - v(0)c(0).$$

This equation can be discretized to give one more linear equation in c'_0 and c_0 in addition to the left boundary condition,

$$Ac_0 + Bc'_0 = G.$$

For details we refer the reader to [5].

3 Algorithm for the inverse problem

The problem we consider is as follows: Given a sampling of the solution $c(x)$ of the system (1) at n various points (possibly with noise) and all parameters except $D(x)$, to recover an approximation to $D(x)$.

Suppose that $c(x)$ solves the system $Lc = f$ together with the boundary conditions. In the setting of an inverse problem, we are usually given the solution $c(x)$ by point-wise measurements d , which may have noise as well. If the operator H represents point-wise evaluation, then the appropriate operator for inversion is

$$F(D) = d = Hc = HL^{-1}(f)$$

and we want to solve $F(D) = d$ for the parameter D . We assume the direct problem is well posed, so that F is well defined. As usual, "inverting" F is

an ill-posed problem, so we resort to the standard Tikhonov regularization technique and cast the problem in the form:

Minimize the functional

$$T_\alpha(D) = \frac{1}{2} \|F(D) - d\|^2 + \alpha \|P(D)\|^2.$$

We now proceed to discretize this operator T_α . The semi-infinite interval $(0, \infty)$ poses special difficulties, since the coordinates of the first and second term of the Tikhonov functional are quite different in nature.

Penalty Term $P(D)$: We choose to define the penalty function in terms of the defining constituents of a solution of the direct problem and a possible nonzero value at infinity. Thus our full model for $D(x)$ (in analogy with the solution $c(x)$ of the direct problem) is

$$D(x) = \tilde{D}(x) + D'(0)q_0(x) + D(0)q_1(x) + D(\infty)q_2(x),$$

where q_0, q_1 are as above, and $q_2(0) = 0, q_2'(0) = 0, q_2(\infty) = 1$. As in Section 2.2, the cardinal functions are chosen in such a way that Theorem 2.1 applies to the reduced function $\tilde{D}(x)$. This amounts to an analyticity restriction on $D(z)$ in the domain \mathcal{D} (which includes the positive x -axis) of Theorem 2.1. Hence we define the penalty function

$$P(D)^2 = a_1 |D'(0)|^2 + a_2 |D(0)|^2 + a_3 |D(\infty)|^2 + a_4 \left\| \tilde{D} \right\|^2 + a_5 \left\| \tilde{D}' \right\|^2.$$

This function discretizes very nicely and Sinc theory can be used to approximate the last two terms exponentially accurately, thanks to our assumptions about D . In practice, the coefficient of $D(\infty)$ tends to be unimportant relative to the others, except for the fact that it tends to overly damp the solution when larger values of the regularization parameter α . Hence, in the absence of additional information, we typically take all coefficients to be 1 except a_3 , which is chosen to be small, say $a_3 = 0.001$.

Functional Term: The term $\frac{1}{2} \|F(D) - d\|^2$ is more problematic since the sampled values

$$d \approx [c(x_k)]_{k=1}^n$$

are not necessarily at the sinc nodes (in fact, almost certainly not), yet in the limit this term should provide global information about D itself. We choose the simplest possible interpretation: that $F(D) - d$ is a fixed vector of values and we simply use the Euclidean norm (scaled by $1/\sqrt{n}$, where n is the dimension of the vector so as to approximate integrals in the limit of n) on this part of the Tikhonov functional.

M_x	$\alpha = 1, \beta = 1, d = \pi/3$			
	m	h	$ e(0) $	$\ e(x_k)\ _\infty$
16	22	0.45	1.75e-03	3.39e-03
32	40	0.32	4.89e-04	5.04e-04
64	78	0.22	2.27e-05	2.78e-05
128	148	0.16	1.69e-08	1.28e-06

Table 1: Accuracy of Sinc approximation for Example 1.

Discretization is now very straightforward: Approximate $L^{-1}(f)$ on a sinc grid, then use a projection matrix H to evaluate the approximate sinc solution at the sampling nodes. Thus, if there are m sinc nodes and n sampling nodes, $L^{-1}(f)$ is discretized to a vector $v = [v_j]_{j=1}^m$ and the discrete version of $F(D) - d$ is

$$F_n(D) - d = Hv - d.$$

4 Examples

We shall illustrate the method by several examples. These examples are somewhat artificial in the sense that the exact answer is known in advance and is even used to generate parameters of the problem. Nonetheless, such an approach is needed to simulate the generation of data and to evaluate the accuracy of the algorithm.

Example 1. The direct algorithm is illustrated by the following example, which is taken from [5]. Consider the steady-state convection-diffusion problem

$$\begin{aligned} -((1 - 0.5e^{-2x})c_x)_x + c_x + c &= f(x), \quad x > 0 \\ c(0) - 0.5c_x(0) &= 1 \\ c(\infty) &= 0. \end{aligned}$$

Here the solution that generated the problem is $c(x) = e^{-x^2/4}$. Some results are tabulated below. Similar results are found over a range of examples. The parameters α, β, d play the same role as Theorem 2.1. For most inverse problems that we have considered, a sinc system of order m between $m = 40$ and $m = 80$ suffices.

Example 2. Consider the direct problem which constitutes Example 1. Suppose that we are given a sampling of the true solution at n equally spaced nodes (equal spacing is not a requirement but a convenience here). Take as

n	BFGS	Error	α	m (10% noise)	BFGS	Error	α
4	95	0.01	0.00048	4	74	0.02	0.00083
20	113	0.0047	0.00022	20	96	0.014	0.001
40	108	0.0047	0.00022	40	90	0.038	0.00083

Table 2: Computations for Example 2.

starting point the constant function $D_0 = 2$ and sinc system size $m = 40$. We sample the true solution both with and without noise at evenly spaced points across the interval $[0.2, 2.4]$. The reason for this choice is that one can see on examination of computed solutions that the solution vanishes very rapidly. In practical measurements one would expect error to swamp very small values; hence it is preferable to sample in regions where the forward solution is clearly nonzero. In the case of noise we assume a relative error of at most 10% due to a uniform distribution.

The number of variables in the Tikhonov functional varied from 87 to 123 in this example. A modified Newton BFGS method with Armijo line search applied to the Tikhonov functional proved to be reasonably effective and avoided the expense of full Hessian calculations. The results of our calculations are displayed in Table 2 for cases of exact data and data with of a uniformly distributed random (relative) error of at most 10%. The number of iterations is listed in the BFGS column. The error displayed is the *direct error*, i.e., the approximate L_2 norm of $c(x) - c_{approx}(x)$, where $c_{approx}(x)$ is the solution to the direct problem obtained by using the computed value of the parameter $D(x)$ in the direct problem and $c(x)$ is the exact solution. Sinc quadrature is used to estimate this norm. Also listed is the choice of regularization parameter suggested by L -curve analysis. Note that $c_{approx}(x)$ is reasonably close to $c(x)$. Of course, closeness of $c_{approx}(x)$ to $c(x)$ is no guarantee that the computed value $D_{approx}(x)$ will be similarly close to the exact diffusion coefficient $D(x)$. Graphs of exact and computed values for the parameter $D(x)$ are presented in Figure 1.

The choice of regularization parameter α is a complex issue which is complicated by the fact that our problem is non-linear. In these examples we used the L -curve approach. However, we use the penalty functional $J(D)$ in place of $\|D\|$ since the latter term is not even defined in the traditional sense, while $J(D)$ serves as a substitute for a norm. It is well known that the L -curve selection method need not be a convergent strategy (see Vogel [7]) even for linear inverse problems. We use the L -curve as a tool to help us achieve a balance between regularization and residual terms. Therefore, irregularities in the behavior of the L -curve are not a matter of special concern. In practice,

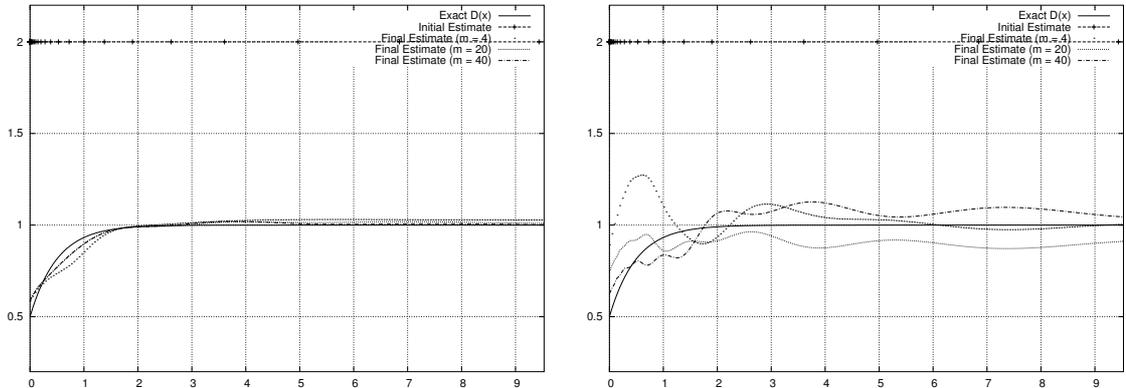


Figure 1: Approximations to parameter $D(x)$ in Example 2.

the L -curve seems to be a very effective tool. For a detailed analysis of various strategies for choice of a regularization parameter see, e. g., Kilmer and O’Leary [3] or the text [8] by Vogel. We conservatively choose the largest α at which a significant corner occurs. In most cases this is graphically obvious, although in a few cases the L -curve was so smooth that a “corner” was difficult to identify. In such cases we approximated the curvature numerically and selected the point of maximum numerical curvature. In a few other cases the L -curve was erratic and several choices for a “corner” were possible.

Example 3. Our final steady-state convection-diffusion problem exhibits singularities in both the coefficients and the solution $c(x) = x^{3/2}e^{-x}$ which is used to generate the right hand side function in the system

$$\begin{aligned} -c_{xx} + c_x + \frac{3}{4x^2}c &= f(x), \quad x > 0 \\ c(0) - c_x(0) &= 0 \\ c(\infty) &= 0. \end{aligned}$$

The singularities make the problem somewhat more intractable than the previous example. In this problem fairly conservative choices of the sinc parameters $\alpha = 1$, $\beta = 1$, $d = \pi/3$ (the same as in Example 1) work reasonably well. Our results are summarized in Table 3 and Figure 2 for the same types of error and measurements as in Example 2.

A natural speculation about these examples is that sinc nodes might yield better results than the equally spaced nodes used in Examples 2 and 3. To compare results we revisit Example 2 and replace equally spaced nodes by sinc nodes. In order for the comparison to be fair, we must sample no more nodes than we do in the equally spaced case. The number of sample points n is set equal to the number of sinc nodes m , which constrains the possibilities

m	BFGS	Error	α	m (10% noise)	BFGS	Error	α
4	93	0.12	0.00022	4	115	0.13	0.00022
20	103	0.0024	0.00022	20	118	0.013	0.0004
40	95	0.0028	0.00022	40	85	0.035	0.00091

Table 3: Computations for Example 3.

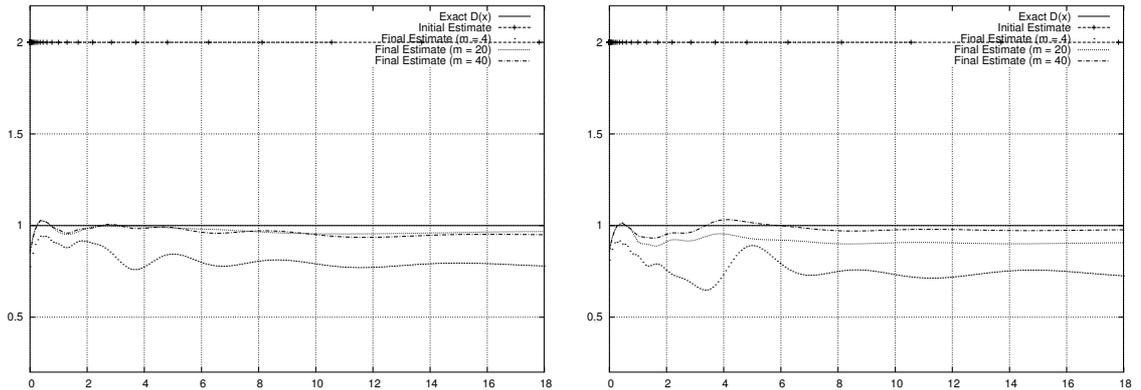


Figure 2: Approximations to parameter $D(x)$ in Example 3.

for n . The choices $n = m = 20$ and $n = m = 40$ can be generated by taking $M_x = 15$ and $M_x = 32$, respectively. A table of the resulting calculations are given in Table 4. A comparison of Tables 3 and 4 shows that sampling at the sinc nodes need not be an optimal strategy.

5 Conclusions

The foregoing algorithm and experiments suggest several conclusions:

1. Sinc methods provide a useful tool for solving certain direct singular problems with high accuracy. These solutions can even exhibit mild derivative singularities at the endpoints of the domain of the forward solution such as $c(x) = x^{1/2}e^{-x}$ or have singular coefficients such as $\lambda(x) = 3/4x^2$.

m	BFGS	Error	α	m (10% noise)	BFGS	Error	α
20	103	0.0095	0.00022	20	108	0.027	0.0005
40	97	0.0138	0.00022	40	111	0.060	0.00083

Table 4: Computations for Example 3 with sampling at sinc nodes.

2. Since high accuracy is obtained with relatively small discretization numbers, minimizers used in solving inverse problems end up working with relatively low dimensional systems.
3. The method is fairly robust with respect to noise, provided that the sinc parameters have been carefully chosen for the forward solver and a suitable choice of regularization parameter α is made.
4. The algorithms we have outlined may prove to be useful in the design of real experiments. For example, we saw that placement of sampling nodes impacts quality of results. Further investigations may suggest optimal placements. Also, the issue of uncertainty in measurement warrants further study. Our examples suggest that more data points in the presence of measurement error may not be better. One possibility is to use fewer nodes with repeated measurements at each node.

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