1. We consider the behavior (for $\lambda \gg 1$) of

$$I(\lambda) = \int_{a}^{b} f(t)e^{i\lambda g(t)} dt \tag{1}$$

where f and g are smooth enough to admit Taylor approximations near some appropriate point in [a, b], and g is real-valued.

2. Suppose that g'(c) = 0 at some point $c \in (a, b)$, and that $g'(t) \neq 0$ everywhere else in the closed interval. Assume moreover that $g''(c) \neq 0$ and $f(c) \neq 0$. Let μ be the sign of g''(c). Thus

$$\mu g''(c) = |g''(c)|.$$

We rewrite $I(\lambda)$ as

$$I(\lambda) = e^{i\lambda g(c)} \int_a^b f(t)e^{i\lambda[g(t) - g(c)]} dt.$$

By the Coates-Euler formula, $\exp(i\lambda [g(t) - g(c)])$ is highly oscillatory for $t \neq c$ and $\lambda \gg 1$. The oscillation gives rise to cancellation which in turn causes the integral to decay rapidly except in a small neighborhood of c. Thus,

$$I(\lambda) \approx e^{i\lambda g(c)} \int_{c-\varepsilon}^{c+\varepsilon} f(t)e^{i\lambda[g(t)-g(c)]} dt$$

$$\approx f(c)e^{i\lambda g(c)} \int_{c-\varepsilon}^{c+\varepsilon} e^{\frac{i\lambda}{2}g''(c)(t-c)^2} dt$$

$$\approx f(c)e^{i\lambda g(c)} \int_{-\infty}^{\infty} e^{\frac{i\lambda}{2}g''(c)(t-c)^2} dt$$

$$= f(c)e^{i\lambda g(c)} \int_{-\infty}^{\infty} e^{\frac{i\lambda}{2}g''(c)s^2} ds$$

$$= f(c)e^{i\lambda g(c)} \sqrt{\frac{2\pi i}{\lambda g''(c)}}$$

$$= f(c)e^{i\lambda g(c)} \sqrt{\frac{2\pi}{\lambda |g''(c)|}} (i\mu)^{\frac{1}{2}}$$

$$= f(c)e^{i\lambda g(c)} \sqrt{\frac{2\pi}{\lambda |g''(c)|}} e^{\frac{\pi i\mu}{4}},$$

for $\lambda \gg 1$. So to leading order

$$I(\lambda) \sim f(c)e^{i\lambda g(c)}\sqrt{\frac{2\pi}{\lambda|g''(c)|}}e^{\frac{\pi i\mu}{4}}, \quad \text{as } \lambda \to \infty.$$
 (2)

Since the main contribution to the integral comes from a region of a point c at which the phase g(t) is stationary, (2) is called the stationary phase approximation.

3. If g(t) is stationary at an endpoint (say t = a) then by the usual modification we obtain the stationary phase approximation

$$I(\lambda) \sim f(a)e^{i\lambda g(a)}\sqrt{\frac{\pi}{2\lambda|g''(a)|}}e^{\frac{\pi i\mu}{4}}, \quad \text{as } \lambda \to \infty,$$
 (3)

where μ is the sign of g''(a).

4. Example: For a fixed integer n, the Bessel function of the first type has the integral representation

$$J_n(\lambda) = \int_0^1 \cos(n\pi t - \lambda \sin \pi t) dt$$
$$= \Re \left\{ \int_0^1 e^{n\pi i t} e^{-i\lambda \sin \pi t} dt \right\}.$$

In the interval [0,1] the phase $g(t) = -\sin \pi t$ is stationary only at the interior point

$$c = \frac{1}{2},$$

with

$$g(c) = -1$$
, $g''(c) = \pi^2$, and $\mu = 1$.

Set

$$f(t) = e^{n\pi i t},$$

so that

$$f(c) = e^{\frac{n\pi i}{2}}.$$

Hence, to leading order,

$$J_n(\lambda) \sim \Re \left\{ e^{\frac{n\pi i}{2}} e^{-i\lambda} \sqrt{\frac{2}{\pi \lambda}} e^{\frac{i\pi}{4}} \right\},$$

$$= \sqrt{\frac{2}{\pi \lambda}} \Re \left\{ e^{-i\left(\lambda - \frac{n\pi}{2} - \frac{\pi}{4}\right)} \right\}$$

$$= \sqrt{\frac{2}{\pi \lambda}} \cos \left(\lambda - \frac{n\pi}{2} - \frac{\pi}{4}\right), \text{ as } \lambda \to \infty.$$