## The Fourier Transform 2

- 1. In class we used eigenfunction expansions to prove the Fourier inversion theorem and Plancherel's identity for "nice" functions of compact support on **R**. You can also use approximate identities to establish these results. Here are some proofs (with a little handwaving).
- **2. Proposition:** If u(x) decays rapidly as  $|x| \to \infty$ , and  $\hat{u}(\xi)$  as  $|\xi| \to \infty$ , then

$$u(x) = \int \hat{u}(\xi)e^{2\pi i\xi \cdot x} d\xi. \tag{1}$$

Proof: Let  $\varepsilon > 0$ . Then

$$\int \hat{u}(\xi)e^{2\pi i\xi \cdot x} d\xi = \int \left[ \int u(y)e^{-2\pi i\xi \cdot y} dy \right] e^{2\pi i\xi \cdot x} d\xi$$

$$= \lim_{\varepsilon \to 0} \int u(y) \left[ \int e^{2\pi i\xi \cdot (x-y)} e^{-4\varepsilon \pi^2 |\xi|^2} d\xi \right] dy$$

$$= \lim_{\varepsilon \to 0} \int G_{\varepsilon}(x-y)u(y) dy$$

$$= \int \delta(x-y)u(y) dy$$

$$= u(x).$$

**3.** Note: You have to take some care with the interpretation of (1). In what sense does equality hold? If u decays rapidly as  $|x| \to \infty$  as is continuous, then equality in (1) is pointwise. If u is in  $L^2(\mathbf{R}^n)$  but not necessarily continuous, then equality holds in the sense of  $L^2$ :

$$\int |u(x) - (\hat{u})(x)|^2 dx = 0.$$

**4. Proposition**: For u and v in  $L^2(\mathbf{R}^n)$ ,

$$\langle \hat{u}, \hat{v} \rangle = \langle u, v \rangle \tag{2}$$

and hence

$$\|\hat{u}\|_2 = \|u\|_2. \tag{3}$$

These are versions of the Plancherel identity.

Proof: We'll use the fact that the Gauss kernel is an approximate identity.

$$\begin{split} \langle \hat{u}, \hat{v} \rangle &= \int \hat{u}(\xi) \bar{\hat{v}}(\xi) \, d\xi \\ &= \int \left[ \int u(x) e^{-2\pi i \xi \cdot x} \, dx \right] \left[ \int \bar{v}(y) e^{-2\pi i \xi \cdot y} \, dy \right] \, d\xi \\ &= \lim_{\varepsilon \to 0} \int u(x) \int \bar{v}(y) \int e^{-4\varepsilon \pi^2 |\xi|^2} e^{2\pi i \xi \cdot (x-y)} \, d\xi \, dy \, dx \\ &= \lim_{\varepsilon \to 0} \int u(x) \int \bar{v}(y) G_{\varepsilon}(y-x) \, dy \, dx \\ &= \int u(x) \int \bar{v}(y) \delta(y-x) \, dy \, dx \\ &= \int u(x) \bar{v}(x) \, dx \\ &= \langle u, v \rangle. \end{split}$$

To get (3) from (2), just take v = u.

5. An operator that preserves inner products is called unitary. Since

$$\langle \mathcal{F}u, \mathcal{F}v \rangle = \langle u, v \rangle$$

the Fourier transform is a unitary operator on  $L^2(\mathbf{R}^n)$ .

**6**. Let  $\varepsilon > 0$  and  $\xi \in \mathbf{R}$ . Find the function  $P_{\varepsilon}(x)$  whose Fourier transform is

$$\hat{P}_{\varepsilon}(\xi) = e^{-2\pi\varepsilon|\xi|}.$$

By the inversion theorem,

$$P_{\varepsilon}(x) = \int \hat{P}_{\varepsilon}(\xi) e^{2\pi i \xi x} d\xi = \frac{1}{\pi} \frac{\varepsilon}{\varepsilon^2 + x^2}.$$

The function  $P_{\varepsilon}$  is called the Possion kernel (for the upper half-plane). It is also an approximate identity as  $\varepsilon \downarrow 0$ .

**7**. The characteristic or indicator of a set A is

$$\chi_A(x) = \begin{cases} 1 & \text{for } x \in A, \\ 0 & \text{for } x \neq A. \end{cases}$$

Let A = [-a, a] in **R**. The function whose Fourier transform is  $\chi_A(\xi)$  is

$$f(x) = \check{\chi}_A(x)$$

$$= \int \chi(\xi) e^{2\pi i \xi x} d\xi$$

$$= \frac{e^{2\pi i a x} - e^{-2\pi i a x}}{2\pi i x}$$

$$= \frac{\sin(2\pi a x)}{\pi x}.$$

8. We can extend the Fourier and inverse Fourier transforms to objects like the Dirac delta function (not really a function) and to functions like  $f(x) \equiv 1$ , for which the Fourier integral (1) doesn't converge. By the defining property of the delta function,

$$\mathcal{F}[\delta(x-y)] = \int \delta(x-y)e^{-2\pi i\xi \cdot x} d\xi = e^{-2\pi i\xi \cdot y}.$$
 (4)

If y = 0, this becomes

$$\hat{\delta}(x) = 1.$$

Thus,

$$\mathcal{F}^{-1}[e^{-2\pi i\xi \cdot y}] = \delta(x - y),$$

and

$$\check{1}(x) = \delta(x).$$

By the definition of the Fourier and inverse Fourier transforms,

$$\hat{f}(\xi) = \check{f}(-\xi). \tag{5}$$

And thus, formally, for fixed  $y \in \mathbf{R}^n$ ,

$$\mathcal{F}[e^{2\pi i\xi \cdot y}] = \mathcal{F}^{-1}[e^{-2\pi i\xi \cdot y}] = \delta(x - y).$$

This sort of formal calculation can be made rigorous with the theory of tempered distributions.